



# Iterative Laplace Transform Method for Solving Fractional Heat and Wave-Like Equations

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## Abstract

In this paper, we derive the closed form solutions of the fractional heat and wave like equations in terms of Mittag-Leffler functions by the use of iterative Laplace transform method. In the process the time-fractional derivatives are considered in Caputo sense for the said problem.

**Keywords:** Laplace transform, Iterative Laplace transform method, heat and wave-like equations, Caputo fractional derivative, Mittag-Leffler function, fractional differential equation. **MSC (2010):** 26A33, 33E12, 35R11, 44A10.

## Introduction

Fractional calculus has been attracting the attention of scientists and engineers from long time ago, resulting in the development of many applications<sup>1-3</sup>. Various methods for the solution of fractional differential equations are available in literature, including fractional subequation method<sup>4</sup>, fractional wavelet method<sup>5-8</sup>, fractional Laplace Adomian decomposition method<sup>9,10</sup>, fractional operational matrix method<sup>11,12</sup>, fractional variational iteration method<sup>13,14</sup>, fractional improved homotopy perturbation method<sup>15,16</sup>, fractional differential transform method<sup>17</sup> and fractional complex transform method<sup>18</sup>.

The iterative method was introduced in 2006 by Daftardar-Gejji and Jafari to solve numerically the nonlinear functional equations<sup>19,20</sup>. By now, the iterative method has been used to solve many integer and fractional boundary value problem<sup>21, 22</sup>. Jafari et al. firstly solved the fractional partial differential equations by the use of iterative Laplace transform method (ILTM)<sup>23</sup>. More recently, Fractional Fokker-Planck equations are solved by the ILTM<sup>24</sup>.

For the present problem, we considered the fractional heat and wave-like equations with variable coefficients in the following form:

$$D_t^\alpha u = f(x, y, z) \frac{\partial^2 u}{\partial x^2} + g(x, y, z) \frac{\partial^2 u}{\partial y^2} + h(x, y, z) \frac{\partial^2 u}{\partial z^2}, \quad (1)$$

and initial conditions:

$$u(x, y, z, 0) = \tilde{h}(x, y, z), u_t(x, y, z, 0) = \ell(x, y, z), \quad (2)$$

where  $\alpha(0 < \alpha \leq 2)$  denotes the fractional derivative. In the case, when  $0 < \alpha \leq 1$ , and  $1 < \alpha \leq 2$ ; then equation. (1) leads to a fractional heat-like and wave-like equations with variable coefficients, respectively.

## Preliminaries and Notations

In this section, we give some basic definitions and properties of fractional calculus and Laplace transform theory, which shall be used in this paper:

**Definition:** The Caputo fractional derivative of function  $u(x, t)$  is defined as<sup>25</sup>

$$D_t^\alpha u(x, t) = \frac{1}{\Gamma(q-\alpha)} \int_0^t (t-v)^{q-\alpha-1} u^{(q)}(x, v) dv, \quad q-1 < \alpha \leq q, \quad q \in N, \quad (3)$$

$$= J_t^{q-\alpha} D^q u(x, t).$$

here  $D^q \equiv \frac{d^q}{dt^q}$  and  $J_t^\alpha$  stands for the Riemann-Liouville

fractional integral operator of order  $\alpha > 0$  defined as<sup>26</sup>

$$J_t^\alpha u(x, t) = \frac{1}{\Gamma(\alpha)} \int_0^t (t-v)^{\alpha-1} u(x, v) dv, \quad v > 0, (q-1 < \alpha \leq q), q \in N \quad (4)$$

**Definition:** The Laplace transform of a function  $\phi(t), t > 0$  is defined as

$$L[\phi(t)] = \Phi(s) = \int_0^\infty e^{-st} \phi(t) dt. \quad (5)$$

**Definition:** Laplace transform of  $D_t^\alpha u(x, t)$  is given as<sup>27</sup>

$$L[D_t^\alpha u(x, t)] = L[u(x, t)] - \sum_{k=0}^{q-1} u^{(k)}(x, 0) s^{\alpha-k-1}, \quad q-1 < \alpha \leq q, \quad q \in N \quad (6)$$

where  $u^{(k)}(x, 0)$  is the k-order derivative of  $u(x, t)$  at  $t = 0$ .

**Definition:** The Mittag-Leffler function which is a generalization of exponential function is defined as<sup>28</sup>

$$E_\alpha(z) = \sum_{q=0}^\infty \frac{z^q}{\Gamma(\alpha q + 1)} (\alpha \in C, \text{Re}(\alpha) > 0). \quad (7)$$

a further generalization of (7) is given in the form <sup>29</sup>

$$E_{\alpha, \beta}(z) = \sum_{q=0}^{\infty} \frac{z^q}{\Gamma(\alpha q + \beta)}; (\alpha, \beta \in C, R(\alpha) > 0, R(\beta) > 0). \quad (8)$$

### Basic Idea of Iterative Laplace Transform Method

To illustrate the basic idea of this method, we consider a general fractional nonlinear nonhomogeneous partial differential equation with the initial conditions of the form:

$$D_t^\alpha u(x,t) + Ru(x,t) + Nu(x,t) = g(x,t), \quad q-1 < \alpha \leq q, \quad q \in N \quad (9)$$

$$u^k(x,0) = h_k(x), \quad k = 0, 1, 2, \dots, q-1 \quad (10)$$

where  $D_t^\alpha u(x,t)$  is the Caputo fractional derivative of the function  $u(x,t)$ ,  $R$  is the linear differential operator,  $N$  represents the general nonlinear differential operator and  $g(x,t)$  is the source term. Applying the Laplace transform (denoted by  $L$  throughout the present paper) in Equation (9), we get

$$L[D_t^\alpha u(x,t)] + L[Ru(x,t) + Nu(x,t)] = L[g(x,t)]. \quad (11)$$

Using Equation (6), we have

$$L[u(x,t)] = \frac{1}{s^\alpha} \sum_{k=0}^{q-1} s^{\alpha-1-k} u^k(x,0) + \frac{1}{s^\alpha} L[g(x,t)] - \frac{1}{s^\alpha} L[Ru(x,t) + Nu(x,t)]. \quad (12)$$

Taking inverse Laplace transform of Equation (12) implies

$$u(x,t) = L^{-1} \left[ \frac{1}{s^\alpha} \left( \sum_{k=0}^{q-1} s^{\alpha-1-k} u^k(x,0) + L[g(x,t)] \right) \right] - L^{-1} \left[ \frac{1}{s^\alpha} L[Ru(x,t) + Nu(x,t)] \right], \quad (13)$$

Now we apply the Iterative method,

$$u(x,t) = \sum_{i=0}^{\infty} u_i(x,t) \quad (14)$$

Since  $R$  is a linear operator,

$$R \left( \sum_{i=0}^{\infty} u_i(x,t) \right) = \sum_{i=0}^{\infty} R(u_i(x,t)) \quad (15)$$

and the nonlinear operator  $N$  is decomposed as

$$N \left( \sum_{i=0}^{\infty} u_i(x,t) \right) = N(u_0(x,t)) + \sum_{i=1}^{\infty} \left\{ N \left( \sum_{k=0}^i u_k(x,t) \right) - N \left( \sum_{k=0}^{i-1} u_k(x,t) \right) \right\} \quad (16)$$

Substituting (14), (15) and (16) in (13), we get

$$\sum_{i=0}^{\infty} u_i(x,t) = L^{-1} \left[ \frac{1}{s^\alpha} \left( \sum_{k=0}^{q-1} s^{\alpha-1-k} u^k(x,0) + L[g(x,t)] \right) \right] - L^{-1} \left[ \frac{1}{s^\alpha} L \left[ \sum_{i=0}^{\infty} R(u_i(x,t)) + N(u_0(x,t)) + \sum_{i=1}^{\infty} \left\{ N \left( \sum_{k=0}^i u_k(x,t) \right) - N \left( \sum_{k=0}^{i-1} u_k(x,t) \right) \right\} \right] \right], \quad (17)$$

We define the recurrence relations as

$$u_0(x,t) = L^{-1} \left[ \frac{1}{s^\alpha} \left( \sum_{k=0}^{q-1} s^{\alpha-1-k} u^k(x,0) + L[g(x,t)] \right) \right]$$

$$u_{q+1}(x,t) = -L^{-1} \left[ \frac{1}{s^\alpha} L \left[ R(u_q(x,t)) - \left\{ N \left( \sum_{k=0}^q u_k(x,t) \right) - N \left( \sum_{k=0}^{q-1} u_k(x,t) \right) \right\} \right] \right], \quad q \geq 1$$

$$u_1(x,t) = -L^{-1} \left[ \frac{1}{s^\alpha} L \left[ R(u_0(x,t)) + N(u_0(x,t)) \right] \right] \quad (18)$$

Therefore the  $Q$ -term approximate solution of (9) - (10) in series form is given by

$$u(x,t) \cong u_0(x,t) + u_1(x,t) + u_2(x,t) + \dots + u_q(x,t), \quad q = 1, 2, \dots \quad (19)$$

### Applications

In this section, the fractional heat and wave-like equations with variable coefficients are solved by ILTM.

**Example:** Consider the following one-dimensional fractional heat-like equation:

$$D_t^\alpha u(x,t) = \frac{1}{2} x^2 \frac{\partial^2 u}{\partial x^2}, \quad 0 < \alpha \leq 1, \quad (20)$$

subject to the initial condition

$$u(x,0) = x^2 \quad (21)$$

Applying the Laplace transform in Equation (20) and making use of (21) we get

$$L[u(x,t)] = \frac{x^2}{s} + \frac{1}{2s^\alpha} x^2 L \left[ \frac{\partial^2 u}{\partial x^2} \right] \quad (22)$$

Taking inverse Laplace transform of Equation (22) implies

$$u(x,t) = x^2 + L^{-1} \left[ \frac{1}{2s^\alpha} x^2 L \left[ \frac{\partial^2 u}{\partial x^2} \right] \right] \quad (23)$$

Now, applying the Iterative method, Substituting (14) - (16) into (23) and applying (18), we obtain the components of the solution as follows:

$$u_0(x,t) = u(x,0) = x^2 \quad (24)$$

$$u_1(x,t) = L^{-1} \left[ \frac{1}{2s^\alpha} x^2 L \left[ \frac{\partial^2 u_0}{\partial x^2} \right] \right]$$

$$= x^2 \frac{t^\alpha}{\Gamma(\alpha + 1)} \quad (25)$$

$$u_2(x,t) = L^{-1} \left[ \frac{1}{2s^\alpha} x^2 L \left[ \frac{\partial^2 (u_0 + u_1)}{\partial x^2} \right] \right] - L^{-1} \left[ \frac{1}{2s^\alpha} x^2 L \left[ \frac{\partial^2 u_0}{\partial x^2} \right] \right]$$

$$= x^2 \left( \frac{t^{2\alpha}}{\Gamma(2\alpha + 1)} + \frac{t^\alpha}{\Gamma(\alpha + 1)} \right) - x^2 \frac{t^\alpha}{\Gamma(\alpha + 1)}$$

$$= x^2 \frac{t^{2\alpha}}{\Gamma(2\alpha + 1)} \quad (26)$$

Therefore, the series form solution is given by

$$u(x, t) = u_0(x, t) + u_1(x, t) + u_2(x, t) + \dots$$

$$u(x, t) = x^2 \left[ 1 + \frac{t^\alpha}{\Gamma(\alpha+1)} + \frac{t^{2\alpha}}{\Gamma(2\alpha+1)} + \dots \right] = x^2 E_\alpha(t^\alpha) \quad (27)$$

where  $E_\alpha(t^\alpha)$  is the Mittag-Leffler function, defined by Equation(7)

**Remark:** Setting  $\alpha=1$ , Equation (20) reduced to one-dimensional heat-like equation:

$$\frac{\partial u}{\partial t} = \frac{1}{2} x^2 \frac{\partial^2 u}{\partial x^2},$$

with solution

$$u(x, t) = x^2 e^t. \quad (28)$$

**Example:** Consider the following Two-dimensional fractional heat-like equation:

$$D_t^\alpha u(x, y, t) = \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2}, \quad 0 < \alpha \leq 1, \quad (29)$$

subject to the initial conditions

$$u(x, y, 0) = \sin x \sin y, \quad (30)$$

Applying the Laplace transform in Equation (29) and making use of (30) we get

$$L[u(x, y, t)] = \frac{\sin x \sin y}{s} + \frac{1}{s^\alpha} L \left[ \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right] \quad (31)$$

Taking inverse Laplace transform of Equation (31) implies

$$u(x, t) = \sin x \sin y + L^{-1} \left[ \frac{1}{s^\alpha} L \left[ \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right] \right] \quad (32)$$

Now, applying the Iterative method,

Substituting (14) - (16) into (32) and applying (18), we obtain the components of the solution as follows:

$$u_0(x, y, t) = u(x, y, 0) = \sin x \sin y \quad (33)$$

$$u_1(x, y, t) = L^{-1} \left[ \frac{1}{s^\alpha} L \left[ \frac{\partial^2 u_0}{\partial x^2} + \frac{\partial^2 u_0}{\partial y^2} \right] \right]$$

$$= -2 \sin x \sin y \frac{t^\alpha}{\Gamma(\alpha+1)} \quad (34)$$

$$u_2(x, y, t) = L^{-1} \left[ \frac{1}{s^\alpha} L \left[ \frac{\partial^2 (u_0 + u_1)}{\partial x^2} + \frac{\partial^2 (u_0 + u_1)}{\partial y^2} \right] \right] - L^{-1} \left[ \frac{1}{s^\alpha} L \left[ \frac{\partial^2 u_0}{\partial x^2} + \frac{\partial^2 u_0}{\partial y^2} \right] \right]$$

$$= \sin x \sin y \left( \frac{(-2)^2 t^{2\alpha}}{\Gamma(2\alpha+1)} - \frac{2t^\alpha}{\Gamma(\alpha+1)} \right) + 2 \sin x \sin y \frac{t^\alpha}{\Gamma(\alpha+1)}$$

$$= (-2)^2 \sin x \sin y \frac{t^{2\alpha}}{\Gamma(2\alpha+1)} \quad (35)$$

Therefore, the series form solution is given by

$$u(x, y, t) = u_0(x, y, t) + u_1(x, y, t) + u_2(x, y, t) + \dots$$

$$u(x, y, t) = \sin x \sin y \left[ 1 + \frac{(-2t^\alpha)}{\Gamma(\alpha+1)} + \frac{(-2t^\alpha)^2}{\Gamma(2\alpha+1)} + \dots \right] = \sin x \sin y [E_\alpha(-2t^\alpha)] \quad (36)$$

**Remark 2.** Setting  $\alpha=1$ , Equation (29) reduced to Two-dimensional heat-like equation:

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2},$$

with solution

$$u(x, y, t) = e^{-2t} \sin x \sin y \quad (37)$$

**Example 3.** Consider the following Three-dimensional fractional heat-like equation:

$$D_t^\alpha u(x, y, z, t) = x^4 y^4 z^4 + \frac{1}{36} \left( x^2 \frac{\partial^2 u}{\partial x^2} + y^2 \frac{\partial^2 u}{\partial y^2} + z^2 \frac{\partial^2 u}{\partial z^2} \right), \quad 0 < \alpha \leq 1, \quad (38)$$

Subject to the initial condition

$$u(x, y, z, 0) = 0, \quad (39)$$

Applying the Laplace transform in Equation (38) and making use of (39) we get

$$L[u(x, y, z, t)] = \frac{1}{s^\alpha} L(x^4 y^4 z^4) + \frac{x^2}{36} L\left(\frac{1}{s^\alpha} u_{xx}\right) + \frac{y^2}{36} L\left(\frac{1}{s^\alpha} u_{yy}\right) + \frac{z^2}{36} L\left(\frac{1}{s^\alpha} u_{zz}\right) \quad (40)$$

Taking inverse Laplace transform of Equation (40) implies

$$u(x, y, z, t) = L^{-1} \left[ \frac{1}{s^\alpha} L(x^4 y^4 z^4) + \frac{x^2}{36} L\left(\frac{1}{s^\alpha} u_{xx}\right) + \frac{y^2}{36} L\left(\frac{1}{s^\alpha} u_{yy}\right) + \frac{z^2}{36} L\left(\frac{1}{s^\alpha} u_{zz}\right) \right] \quad (41)$$

Now we apply the Iterative method,

Substituting (14) - (16) into (41) and applying (18), we obtain the components of the solution as follows:

$$u_0(x, y, z, t) = x^4 y^4 z^4 \frac{t^\alpha}{\Gamma(\alpha+1)} \quad (42)$$

$$u_1(x, y, z, t) = L^{-1} \left[ \frac{1}{36s^\alpha} x^2 L\left(\frac{\partial^2 u_0}{\partial x^2}\right) + \frac{1}{36s^\alpha} y^2 L\left(\frac{\partial^2 u_0}{\partial y^2}\right) + \frac{1}{36s^\alpha} z^2 L\left(\frac{\partial^2 u_0}{\partial z^2}\right) \right]$$

$$= x^4 y^4 z^4 \frac{t^{2\alpha}}{\Gamma(2\alpha+1)} \quad (43)$$

$$u_2(x, y, z, t) = L^{-1} \left[ \frac{1}{36s^\alpha} x^2 L\left(\frac{\partial^2 (u_0 + u_1)}{\partial x^2}\right) + \frac{1}{36s^\alpha} y^2 L\left(\frac{\partial^2 (u_0 + u_1)}{\partial y^2}\right) + \frac{1}{36s^\alpha} z^2 L\left(\frac{\partial^2 (u_0 + u_1)}{\partial z^2}\right) \right]$$

$$- L^{-1} \left[ \frac{1}{36s^\alpha} x^2 L\left(\frac{\partial^2 u_0}{\partial x^2}\right) + \frac{1}{36s^\alpha} y^2 L\left(\frac{\partial^2 u_0}{\partial y^2}\right) + \frac{1}{36s^\alpha} z^2 L\left(\frac{\partial^2 u_0}{\partial z^2}\right) \right]$$

$$= \left( x^4 y^4 z^4 \frac{t^{2\alpha}}{\Gamma(2\alpha+1)} + x^4 y^4 z^4 \frac{t^{3\alpha}}{\Gamma(3\alpha+1)} \right) - \left( x^4 y^4 z^4 \frac{t^{2\alpha}}{\Gamma(2\alpha+1)} \right)$$

$$= x^4 y^4 z^4 \left( \frac{t^{3\alpha}}{\Gamma(3\alpha+1)} \right) \quad (44)$$

Therefore, the series form solution is given by  
 $u(x, y, z, t) = u_0(x, y, z, t) + u_1(x, y, z, t) + u_2(x, y, z, t) + \dots$

$$u(x, y, z, t) = x^4 y^4 z^4 \left[ \frac{t^\alpha}{\Gamma(\alpha+1)} + \frac{t^{2\alpha}}{\Gamma(2\alpha+1)} + \frac{t^{3\alpha}}{\Gamma(3\alpha+1)} + \dots \right] = x^4 y^4 z^4 [E_\alpha(t^\alpha) - 1]. \quad (45)$$

**Remark:** Setting  $\alpha=1$ , Equation (38) reduced to Three-dimensional heat-like equation:

$$\frac{\partial u}{\partial t} = x^4 y^4 z^4 + \frac{1}{36} \left( x^2 \frac{\partial^2 u}{\partial x^2} + y^2 \frac{\partial^2 u}{\partial y^2} + z^2 \frac{\partial^2 u}{\partial z^2} \right),$$

with solution

$$u(x, y, z, t) = x^4 y^4 z^4 (e^t - 1) \quad (46)$$

**Example:** Consider the following one-dimensional fractional wave-like equation:

$$D_t^\alpha u(x, t) = \frac{1}{2} x^2 \frac{\partial^2 u}{\partial x^2}, \quad 1 < \alpha \leq 2, \quad (47)$$

subject to the initial conditions

$$u(x, 0) = x, \quad u_t(x, 0) = x^2 \quad (48)$$

Applying the Laplace transform in Equation (47) and making use of (48) we get

$$L[u(x, t)] = \frac{x}{s} + \frac{x^2}{s^2} + \frac{x^2}{2s^\alpha} L\left[\frac{\partial^2 u}{\partial x^2}\right] \quad (49)$$

Taking inverse Laplace transform of Equation (49) implies

$$u(x, t) = x + x^2 t + L^{-1} \left[ \frac{1}{2s^\alpha} L\left[\frac{\partial^2 u}{\partial x^2}\right] \right] \quad (50)$$

Now, applying the Iterative method,

Substituting (14) - (16) into (50) and applying (18), we obtain the components of the solution as follows:

$$u_0(x, t) = x + x^2 t \quad (51)$$

$$u_1(x, t) = L^{-1} \left[ \frac{1}{2s^\alpha} x^2 L\left[\frac{\partial^2 u_0}{\partial x^2}\right] \right] = x^2 \frac{t^{\alpha+1}}{\Gamma(\alpha+2)} \quad (52)$$

$$u_2(x, t) = L^{-1} \left[ \frac{1}{2s^\alpha} x^2 L\left[\frac{\partial^2 (u_0 + u_1)}{\partial x^2}\right] \right] - L^{-1} \left[ \frac{1}{2s^\alpha} x^2 L\left[\frac{\partial^2 u_0}{\partial x^2}\right] \right] \\ = x^2 \left( \frac{t^{2\alpha+1}}{\Gamma(2\alpha+2)} + \frac{t^{\alpha+1}}{\Gamma(\alpha+2)} \right) - \left( x^2 \frac{t^{\alpha+1}}{\Gamma(\alpha+2)} \right) = x^2 \frac{t^{2\alpha+1}}{\Gamma(2\alpha+2)} \quad (53)$$

Therefore, the series form solution is given by

$$u(x, t) = u_0(x, t) + u_1(x, t) + u_2(x, t) + \dots$$

$$u(x, t) = x + x^2 \left[ t + \frac{t^{\alpha+1}}{\Gamma(\alpha+2)} + \frac{t^{2\alpha+1}}{\Gamma(2\alpha+2)} + \dots \right] = x + x^2 t E_{\alpha, 2}(t^\alpha) \quad (54)$$

where  $E_{\alpha, \beta}(t^\alpha)$  is generalization form of Mittag-Leffler function, defined by Equation(8).

**Remark 4.** Setting  $\alpha=2$ , Equation (47) reduced to one-dimensional wave-like equation:

$$\frac{\partial u}{\partial t} = \frac{1}{2} x^2 \frac{\partial^2 u}{\partial x^2},$$

with solution

$$u(x, t) = x + x^2 \sinh t \quad (55)$$

**Example 5.** Consider the following Two-dimensional fractional wave-like equation:

$$D_t^\alpha u(x, y, t) = \frac{1}{12} \left( x^2 \frac{\partial^2 u}{\partial x^2} + y^2 \frac{\partial^2 u}{\partial y^2} \right), \quad 1 < \alpha \leq 2, \quad (56)$$

Subject to the initial conditions

$$u(x, y, 0) = x^4, \quad u_t(x, y, 0) = y^4 \quad (57)$$

Applying the Laplace transform in Equation (56) and making use of (57) we get

$$L[u(x, y, t)] = \frac{x^4}{s} + \frac{y^4}{s^2} + \frac{x^2}{12s^\alpha} L\left[\frac{\partial^2 u}{\partial x^2}\right] + \frac{y^2}{12s^\alpha} L\left[\frac{\partial^2 u}{\partial y^2}\right] \quad (58)$$

Taking inverse Laplace transform of Equation (58) implies

$$u(x, y, t) = x^4 + y^4 t + L^{-1} \left[ \frac{x^2}{12s^\alpha} L\left[\frac{\partial^2 u}{\partial x^2}\right] + \frac{y^2}{12s^\alpha} L\left[\frac{\partial^2 u}{\partial y^2}\right] \right] \quad (59)$$

Now, applying the Iterative method,

Substituting (14) - (16) into (59) and applying (18), we obtain the components of the solution as follows:

$$u_0(x, y, t) = x^4 + y^4 t \quad (60)$$

$$u_1(x, y, t) = L^{-1} \left[ \frac{x^2}{12s^\alpha} L\left[\frac{\partial^2 u_0}{\partial x^2}\right] + \frac{y^2}{12s^\alpha} L\left[\frac{\partial^2 u_0}{\partial y^2}\right] \right] \\ = x^4 \frac{t^\alpha}{\Gamma(\alpha+1)} + y^4 \frac{t^{\alpha+1}}{\Gamma(\alpha+2)} \quad (61)$$

$$u_2(x, y, t) = L^{-1} \left[ \frac{x^2}{12s^\alpha} L\left[\frac{\partial^2 (u_0 + u_1)}{\partial x^2}\right] + \frac{y^2}{12s^\alpha} L\left[\frac{\partial^2 (u_0 + u_1)}{\partial y^2}\right] \right] \\ - L^{-1} \left[ \frac{x^2}{12s^\alpha} L\left[\frac{\partial^2 u_0}{\partial x^2}\right] + \frac{y^2}{12s^\alpha} L\left[\frac{\partial^2 u_0}{\partial y^2}\right] \right] \quad (62)$$

$$= x^4 \left( \frac{t^\alpha}{\Gamma(\alpha+1)} + \frac{t^{2\alpha}}{\Gamma(2\alpha+1)} \right) + y^4 \left( \frac{t^{\alpha+1}}{\Gamma(\alpha+2)} + \frac{t^{2\alpha+1}}{\Gamma(2\alpha+2)} \right) - x^4 \left( \frac{t^\alpha}{\Gamma(\alpha+1)} \right) - y^4 \left( \frac{t^{\alpha+1}}{\Gamma(\alpha+2)} \right) \\ = x^4 \frac{t^{2\alpha}}{\Gamma(2\alpha+1)} + y^4 \frac{t^{2\alpha+1}}{\Gamma(2\alpha+2)}, \quad (63)$$

Therefore, the series form solution is given by

$$u(x, y, t) = u_0(x, y, t) + u_1(x, y, t) + u_2(x, y, t) + \dots$$

$$u(x,y,t) = x^4 \left[ 1 + \frac{t^\alpha}{\Gamma(\alpha+1)} + \frac{t^{2\alpha}}{\Gamma(2\alpha+1)} + \dots \right] + y^4 \left[ 1 + \frac{t^{\alpha+1}}{\Gamma(\alpha+2)} + \frac{t^{2\alpha+1}}{\Gamma(2\alpha+2)} + \dots \right]$$

$$= x^4 E_\alpha(t^\alpha) + y^4 t E_{\alpha,2}(t^\alpha). \quad (64)$$

**Remark 5** Setting  $\alpha=2$ , Equation (56) reduced to Two-dimensional wave-like equation:

$$\frac{\partial u}{\partial t} = \frac{1}{12} \left( x^2 \frac{\partial^2 u}{\partial x^2} + y^2 \frac{\partial^2 u}{\partial y^2} \right),$$

with solution

$$u(x, y, t) = x^4 \cosh t + y^4 \sinh t. \quad (65)$$

**Example 6.** Consider the following Three-dimensional fractional wave-like equation:

$$D_t^\alpha u(x, y, z, t) = x^2 + y^2 + z^2 + \frac{1}{2} \left( x^2 \frac{\partial^2 u}{\partial x^2} + y^2 \frac{\partial^2 u}{\partial y^2} + z^2 \frac{\partial^2 u}{\partial z^2} \right), \quad 1 < \alpha \leq 2, \quad (66)$$

subject to the initial conditions

$$u(x, y, z, 0) = 0, \quad u_t(x, y, z, 0) = x^2 + y^2 - z^2. \quad (67)$$

Applying the Laplace transform in Equation (66) and making use of (67) we get

$$L[u(x, y, z, t)] = \left( \frac{x^2 + y^2 - z^2}{s^\alpha} \right) + \left( \frac{x^2 + y^2 + z^2}{s^\alpha} \right) + \frac{x^2}{2s^\alpha} L \left( \frac{\partial^2 u}{\partial x^2} \right) + \frac{y^2}{2s^\alpha} L \left( \frac{\partial^2 u}{\partial y^2} \right) + \frac{z^2}{2s^\alpha} L \left( \frac{\partial^2 u}{\partial z^2} \right) \quad (68)$$

Taking inverse Laplace transform of Equation (68) implies

$$u(x, y, z, t) = t(x^2 + y^2 - z^2) + L^{-1} \left[ \frac{1}{s^\alpha} L(x^2 + y^2 + z^2) + \frac{x^2}{2s^\alpha} L \left( \frac{\partial^2 u}{\partial x^2} \right) + \frac{y^2}{2s^\alpha} L \left( \frac{\partial^2 u}{\partial y^2} \right) + \frac{z^2}{2s^\alpha} L \left( \frac{\partial^2 u}{\partial z^2} \right) \right] \quad (69)$$

Now, applying the Iterative method,

Substituting (14) - (16) into (69) and applying (18), we obtain the components of the solution as follows:

$$u_0(x, y, z, t) = t(x^2 + y^2 - z^2) + L^{-1} \left[ \frac{1}{s^\alpha} L(x^2 + y^2 + z^2) \right]$$

$$= t(x^2 + y^2 - z^2) + (x^2 + y^2 + z^2) \frac{t^\alpha}{\Gamma(\alpha+1)} \quad (70)$$

$$u_1(x, y, z, t) = L^{-1} \left[ \frac{x^2}{2s^\alpha} L \left( \frac{\partial^2 u_0}{\partial x^2} \right) + \frac{y^2}{2s^\alpha} L \left( \frac{\partial^2 u_0}{\partial y^2} \right) + \frac{z^2}{2s^\alpha} L \left( \frac{\partial^2 u_0}{\partial z^2} \right) \right]$$

$$= (x^2 + y^2 - z^2) \frac{t^{\alpha+1}}{\Gamma(\alpha+2)} + (x^2 + y^2 + z^2) \frac{t^{2\alpha}}{\Gamma(2\alpha+1)} \quad (71)$$

$$u_2(x, y, z, t) = L^{-1} \left[ \frac{x^2}{2s^\alpha} L \left( \frac{\partial^2 (u_0 + u_1)}{\partial x^2} \right) + \frac{y^2}{2s^\alpha} L \left( \frac{\partial^2 (u_0 + u_1)}{\partial y^2} \right) + \frac{z^2}{2s^\alpha} L \left( \frac{\partial^2 (u_0 + u_1)}{\partial z^2} \right) \right]$$

$$- L^{-1} \left[ \frac{x^2}{2s^\alpha} L \left( \frac{\partial^2 u_0}{\partial x^2} \right) + \frac{y^2}{2s^\alpha} L \left( \frac{\partial^2 u_0}{\partial y^2} \right) + \frac{z^2}{2s^\alpha} L \left( \frac{\partial^2 u_0}{\partial z^2} \right) \right]$$

$$= (x^2 + y^2 - z^2) \frac{t^{2\alpha+1}}{\Gamma(2\alpha+2)} + (x^2 + y^2 + z^2) \frac{t^{3\alpha}}{\Gamma(3\alpha+1)} \quad (72)$$

Therefore, the series form solution is given by

$$u(x, y, z, t) = u_0(x, y, z, t) + u_1(x, y, z, t) + u_2(x, y, z, t) + \dots$$

$$u(x, y, z, t) = (x^2 + y^2 - z^2) \left[ t + \frac{t^{\alpha+1}}{\Gamma(\alpha+2)} + \frac{t^{2\alpha+1}}{\Gamma(2\alpha+2)} + \dots \right] + (x^2 + y^2 + z^2) \left[ \frac{t^\alpha}{\Gamma(\alpha+1)} + \frac{t^{2\alpha}}{\Gamma(2\alpha+1)} + \frac{t^{3\alpha}}{\Gamma(3\alpha+1)} + \dots \right]$$

$$= t(x^2 + y^2 - z^2) E_{\alpha,2}(t^\alpha) + (x^2 + y^2 + z^2) [E_\alpha(t^\alpha) - 1]. \quad (73)$$

**Remark 6.** Setting  $\alpha=2$ , Equation (66) reduced to Three-dimensional wave-like equation:

$$\frac{\partial u}{\partial t} = x^2 + y^2 + z^2 + \frac{1}{2} \left( x^2 \frac{\partial^2 u}{\partial x^2} + y^2 \frac{\partial^2 u}{\partial y^2} + z^2 \frac{\partial^2 u}{\partial z^2} \right),$$

with solution

$$u(x, y, z, t) = (x^2 + y^2) e^t + z^2 e^{-t} + (x^2 + y^2 - z^2). \quad (74)$$

## Conclusion

The solutions of the fractional heat and wave like equations in terms of Mittag-Leffler functions by the use of iterative Laplace transform method were derived. The solutions are obtained in series form that rapidly converges in a closed exact formula with simply computable terms. The calculations are simple and straightforward. The method was tested on six examples on different situations.

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